

ADCIRC v55

MODELING THE EARTH, MESH RESOLUTION EFFECTS, AND REMOVING TIME STEP CONSTRAINTS

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OVERVIEW OF COASTAL OCEAN MODELS

O.B. Fringer, C.N. Dawson, R. He et al.

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Table 1
Models mentioned in this paper, in alphabetical order by model name. In the grid/variable placement column, B and C refer to the grid types of [Arakawa and Lamb \(1977\)](#).

Model	Citation	C: Coastal, R: Regional, G: Global	Finite-volume (FV) or Finite-element (FE)	Grid/variable placement	Vertical coordinate	Developer (s) present?	Notes/unique features
ADCIRC	Luettich et al. (1992) Westerink et al. (1994)	C	FE		Sigma	Blaine, Luettich, westerink	Continuous/Discontinuous Galerkin
COAWST	Warner et al. (2008, 2010)	C/R	FV	C	Sigma	Signell, He, Ganju	Coupled Atmo-sphere/wave/sediment
COSINE	Chai et al. (2002, 2003, 2007)	C/R	–	–	–	Chai	Biogeochemical model
Delft3D-Flow/Delft3D-FM	oss.deltares.nl/web/delft3d/home	C	FV	C	Sigma	No	
ECOM-si/POM	Blumberg and Mellor (1987)	C/R	FV	C	Sigma	No	
FunwaveC	Feddersen et al. (2011)	C	FV	C	2D (x-y)	Kirby, Shi	2D Boussinesq wave model
FVCOM	Chen et al. (2003)	C/R/G	FV	B	Sigma	Chen	
GETM	Burchard and Bolding (2002)	C	FV	C	Sigma/Adaptive	Burchard	Numerical mixing analysis
GOTM	Burchard et al. (1999)	C	–	–	1D (Z)	Burchard	Turbulence Model
HYCOM	Bleck and Boudra (1981)	R/G	FV	C	Sigma/Z/Isopycnal	No	Hybrid/Isopycnal coordinates
MARS3D	Lazure and Dumas (2008)	C/R	FV	C	Sigma	No	
MITgcm	Marshall et al. (1997a)	C/R/G	FV	C	Z	No	Shaved cells
NHWAVE	Ma et al. (2012) Shi et al. (2015)	C	FV	C	Sigma	Kirby, Shi	Nonhydrostatic 3D wave model, LES
ROMS	Shchepetkin and McWilliams (2005)	R	FV	C	Sigma	Wilkin	
SCHISM	Zhang et al. (2016)	C/R	FV/FE	C/FE	Z/Sigma	Zhang	Locally-adaptive vertical coordinate
SELF	Zhang and Baptista (2008)	C	FV/FE	C	Z/Sigma	Zhang, Baptista	
SLIM	Vallaeys et al. (2018)	C/R	FE		Z/Sigma/Adaptive	No	Discontinuous Galerkin
SUNTANS	Fringer et al. (2006)	C	FV	C	Z	Fringer	
SWAN	Booij et al. (1999)	C/R	FV	C	2D (x-y)	No	Phase-averaged wave model
SWASH	Zijlema et al. (2011)	C	FV	C	Sigma	No	2D/3D Boussinesq wave model
TRIM/UnTRIM	Casulli (1999) , Casulli and Zanolli (2002, 2005)	C	FV	C	Z	No	Subgrid bathymetry
WaveWatch III	Tolman (2009)	R/G	FV	C	2D (x-y)	No	Phase-averaged wave model
WRF-Hydro	Gochis et al. (2018)	C/R	–	–	–	No	Atmospheric/hydrological model

- **ADCIRC** is considered only a **coastal** model
- Continuous Galerkin Finite-Element method (**CG-FEM**)
- Uses **GWCE** formulation to avoid high-frequency oscillations in **CG-FEM** method

AIMS

- 1) Extend **ADCIRC** to a **global** model
 - a. Show improvement over old version
 - b. Mesh resolution sensitivities
 - 2) Improve **stability and mass-conservation** issues associated with **GWCE** formulation
- Note: **CG-FEM** method is nice because we can keep **2nd order** accuracy and have a lot of **tolerance for skewed elements** (Fringer et al., 2019)

O.B. Fringer, C.N. Dawson, R. He et al.

Table 1

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Model	Citation	C: Coastal, R: Regional, G: Global	Finite-volume (FV) or Finite-element (FE)
ADCIRC	Luetlich et al. (1992) Westerink et al. (1994)	C/R/G	FE
COAWST	Warner et al. (2008, 2010)	C/R	FV
COSINE	Chai et al. (2002, 2003, 2007)	C/R	–
Delft3D-Flow/Delft3D-FM	oss.deltares.nl/web/delft3d/home	C	FV
ECOM-si/POM	Blumberg and Mellor (1987)	C/R	FV
FunwaveC	Feddersen et al. (2011)	C	FV
FVCOM	Chen et al. (2003)	C/R/G	FV
GETM	Burchard and Bolding (2002)	C	FV
GOTM	Burchard et al. (1999)	C	–
HYCOM	Bleck and Boudra (1981)	R/G	FV
MARS3D	Lazure and Dumas (2008)	C/R	FV
MITgcm	Marshall et al. (1997a)	C/R/G	FV
NHWAVE	Ma et al. (2012) Shi et al. (2015)	C	FV
ROMS	Shchepetkin and McWilliams (2005)	R	FV
SCHISM	Zhang et al. (2016)	C/R	FV/FE
SELFE	Zhang and Baptista (2008)	C	FV/FE
SLIM	Vallaey et al. (2018)	C/R	FE
SUNTANS	Fringer et al. (2006)	C	FV
SWAN	Booij et al. (1999)	C/R	FV
SWASH	Zijlema et al. (2011)	C	FV
TRIM/UnTRIM	Casulli (1999), Casulli and Zanolli (2002, 2005)	C	FV
WaveWatch III	Tolman (2009)	R/G	FV
WRF-Hydro	Gochis et al. (2018)	C/R	–

1) EXTENDING ADCIRC TO GLOBAL MODEL

Currently, $\tan(\phi)$ terms ignored...

$$\frac{1}{R \cos \phi} \frac{\partial(VH \cos \phi)}{\partial \phi} = \frac{1}{R} \frac{\partial(VH)}{\partial \phi} - \cancel{\frac{\tan \phi}{R} VH} \quad \text{0}$$

$$\frac{\partial \zeta}{\partial t} = - \frac{1}{R \cos \phi} \left[\frac{\partial(UH)}{\partial \lambda} + \frac{\partial(VH \cos \phi)}{\partial \phi} \right] \quad (1)$$

$$\begin{aligned} \frac{\partial U}{\partial t} = & - \frac{U}{R \cos \phi} \frac{\partial U}{\partial \lambda} - \frac{V}{R} \frac{\partial U}{\partial \phi} - (\mathcal{C}_{\lambda\phi} - f')V - \frac{1}{R \cos \phi} \frac{\partial \Psi}{\partial \lambda} + \tau_w U_w - (\tau_b + \mathcal{C}_{\lambda\lambda})U \\ & + \frac{\nu_t}{R} \left[\frac{1}{\cos \phi} \frac{\partial \tau_{\lambda\lambda}}{\partial \lambda} + \frac{\partial \tau_{\lambda\phi}}{\partial \phi} - \tan \phi (\tau_{\lambda\phi} + \tau_{\phi\lambda}) \right] \end{aligned} \quad (2)$$

$$\begin{aligned} \frac{\partial V}{\partial t} = & - \frac{U}{R \cos \phi} \frac{\partial V}{\partial \lambda} - \frac{V}{R} \frac{\partial V}{\partial \phi} - (\mathcal{C}_{\phi\lambda} + f')U - \frac{1}{R} \frac{\partial \Psi}{\partial \phi} + \tau_w V_w - (\tau_b + \mathcal{C}_{\phi\phi})V \\ & + \frac{\nu_t}{R} \left[\frac{\partial \tau_{\phi\phi}}{\partial \phi} + \frac{1}{\cos \phi} \frac{\partial \tau_{\phi\lambda}}{\partial \lambda} + \tan \phi (\tau_{\lambda\lambda} - \tau_{\phi\phi}) \right] \end{aligned} \quad (3)$$

$$f' = 2\Omega \sin \phi + \cancel{\frac{\tan \phi}{R} U} \quad \text{0}$$

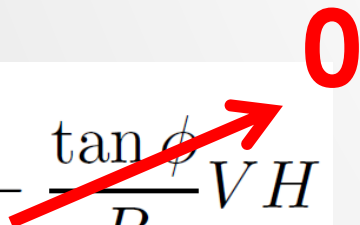
CURRENT ADCIRC MODEL EQUATIONS

- Main problem:

$$\frac{1}{R \cos \phi} \frac{\partial(V H \cos \phi)}{\partial \phi}$$

Solving this term in continuity eq. directly is difficult with CG-FEM due to nonlinearity of the ϕ dependent terms..

$$= \frac{1}{R} \frac{\partial(V H)}{\partial \phi} - \frac{\tan \phi}{R} V H$$



- This expansion eliminates nonlinearity but the $\tan(\phi)$ term is extremely stiff for the numerical method so it has just been ignored...

SOLUTION BY REFORMULATION

- Use an arbitrary cylindrical projection to map (λ, ϕ) onto (x, y) :
(Select desired $p = 0, 1, 2$)

$$x = R(\lambda - \lambda_0) \cos \phi_0$$

$$y = \begin{cases} R \sin \phi \sec \phi_0 \\ R \phi \\ R \ln (\tan \phi + \sec \phi) \cos \phi_0 \end{cases}$$

(λ_0, ϕ_0) is arbitrary origin

if $p = 0$: Equal-area

if $p = 1$: Equidistant (CPP)

if $p = 2$: Conformal (Mercator)

ICS = 2
sans $\tan(\phi)$
terms

ICS = 20

ICS = 21

ICS = 22

- Multiply continuity by $\cos^p(\phi)$ [= 1 when $p = 0$]:

$$\frac{\partial(\zeta \cos^p \phi)}{\partial t} = -L_x \frac{\partial(UH)}{\partial x} - L_y \frac{\partial(VH \cos \phi)}{\partial y}$$

Continuity in a
nice form to solve!

$$L_x = \cos \phi_0 (\cos \phi)^{p-1}, \quad L_y = (\cos \phi_0)^{p-1},$$

this is just a
constant

EQUATIONS ADCIRC NOW SOLVES

$$\frac{\partial^2 (\zeta \cos^p \phi)}{\partial t^2} + \tau_0 \frac{\partial (\zeta \cos^p \phi)}{\partial t} = -L_x \frac{\partial J_x}{\partial x} - L_y \frac{\partial (J_y \cos \phi)}{\partial y}$$

$$J_x = H \left(\frac{\partial U}{\partial t} \right) + U \frac{\partial \zeta}{\partial t} + \tau_0 U H$$

$$J_y = H \left(\frac{\partial V}{\partial t} \right) + V \frac{\partial \zeta}{\partial t} + \tau_0 V H$$

GWCE

τ_0 : just a constant here

$$\frac{\partial U}{\partial t} = -KU \frac{\partial U}{\partial x} - K^{p-1} V \frac{\partial U}{\partial y} - (\mathcal{C}_{xy} - f')V - K \frac{\partial \Psi}{\partial x} + \tau_w U_w - (\tau_b + \mathcal{C}_{xx})U$$

$$+ \nu_t \left[K \frac{\partial \tau_{xx}}{\partial x} + K^{p-1} \frac{\partial \tau_{xy}}{\partial y} - \tan \phi (\tau_{xy} + \tau_{yx}) \right]$$

Momentum

$$\frac{\partial V}{\partial t} = -KU \frac{\partial V}{\partial x} - K^{p-1} V \frac{\partial V}{\partial y} - (\mathcal{C}_{yx} + f')U - K^{p-1} \frac{\partial \Psi}{\partial y} + \tau_w V_w - (\tau_b + \mathcal{C}_{yy})V$$

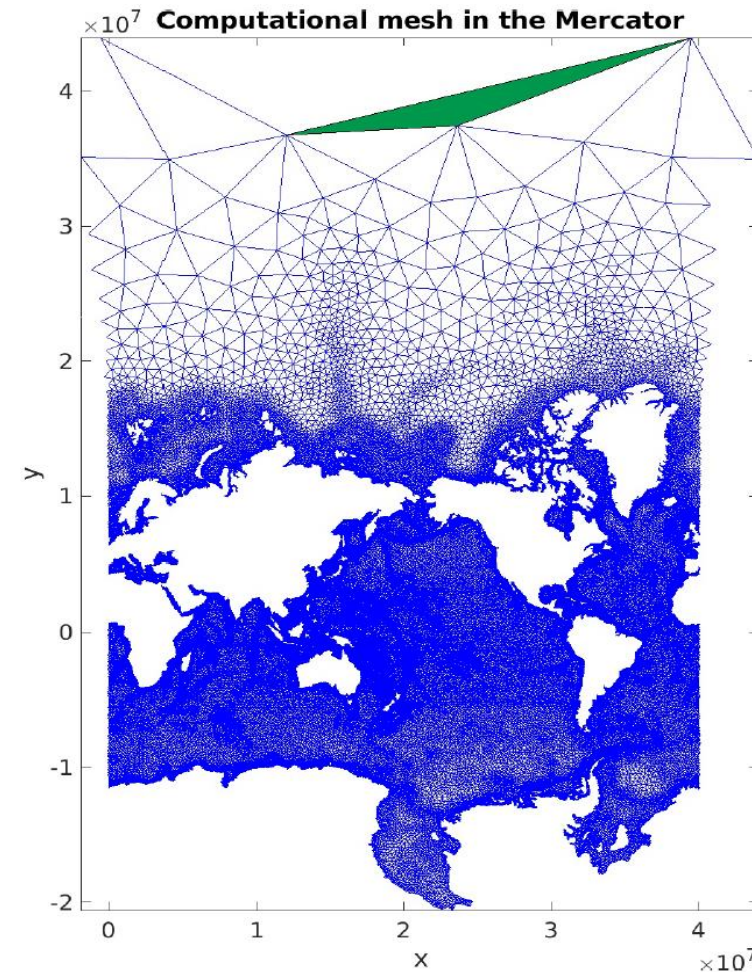
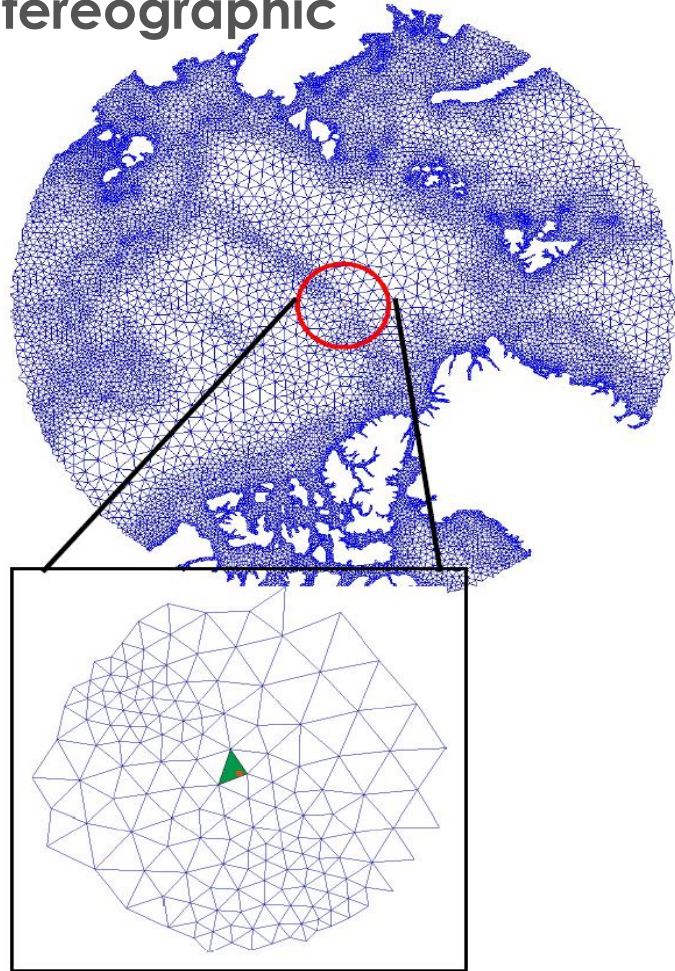
$$+ \nu_t \left[K^{p-1} \frac{\partial \tau_{yy}}{\partial y} + K \frac{\partial \tau_{yx}}{\partial x} + \tan \phi (\tau_{xx} - \tau_{yy}) \right]$$

POLE TREATMENT

- 👉 Strategy 1: do not nothing; requires that no FE node is placed at the pole.

MESH IN ARCTIC REGION

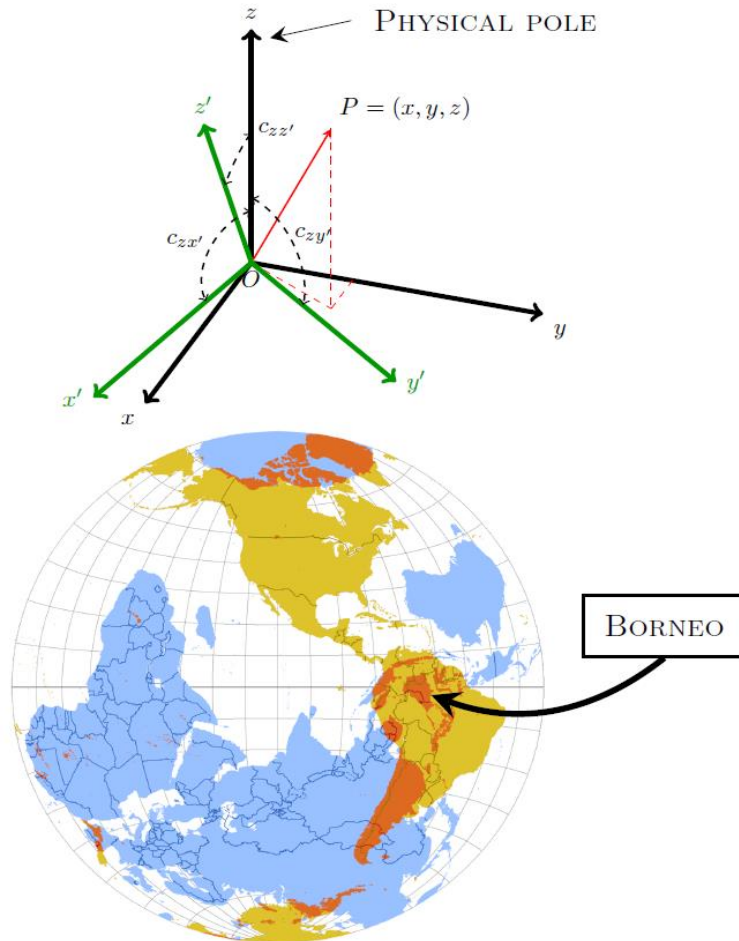
stereographic



- Performs satisfactory but eventually suffers from numerical instability when the element straddling the pole having nodes very close to the pole.

POLE TREATMENT

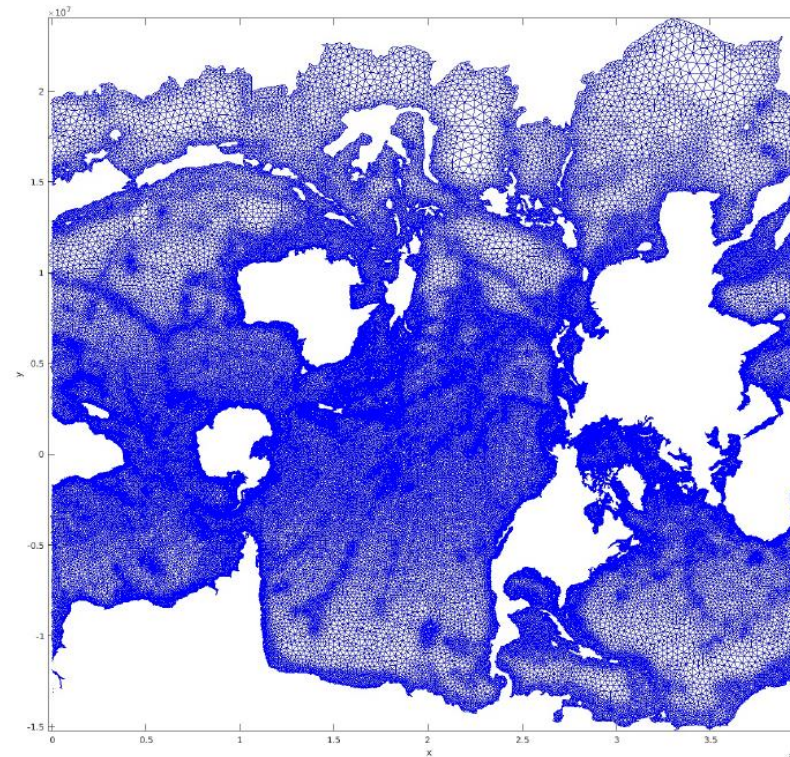
👉 Strategy 2: rotate a computational pole so that it goes through lands.



ANTIPOLES MAP (FROM WIKIPEDIA)

COMPUTATIONAL POLE: BORNEO-BRAZIL

COMPUTATIONAL MESH IN THE MERCATOR PROJECTION

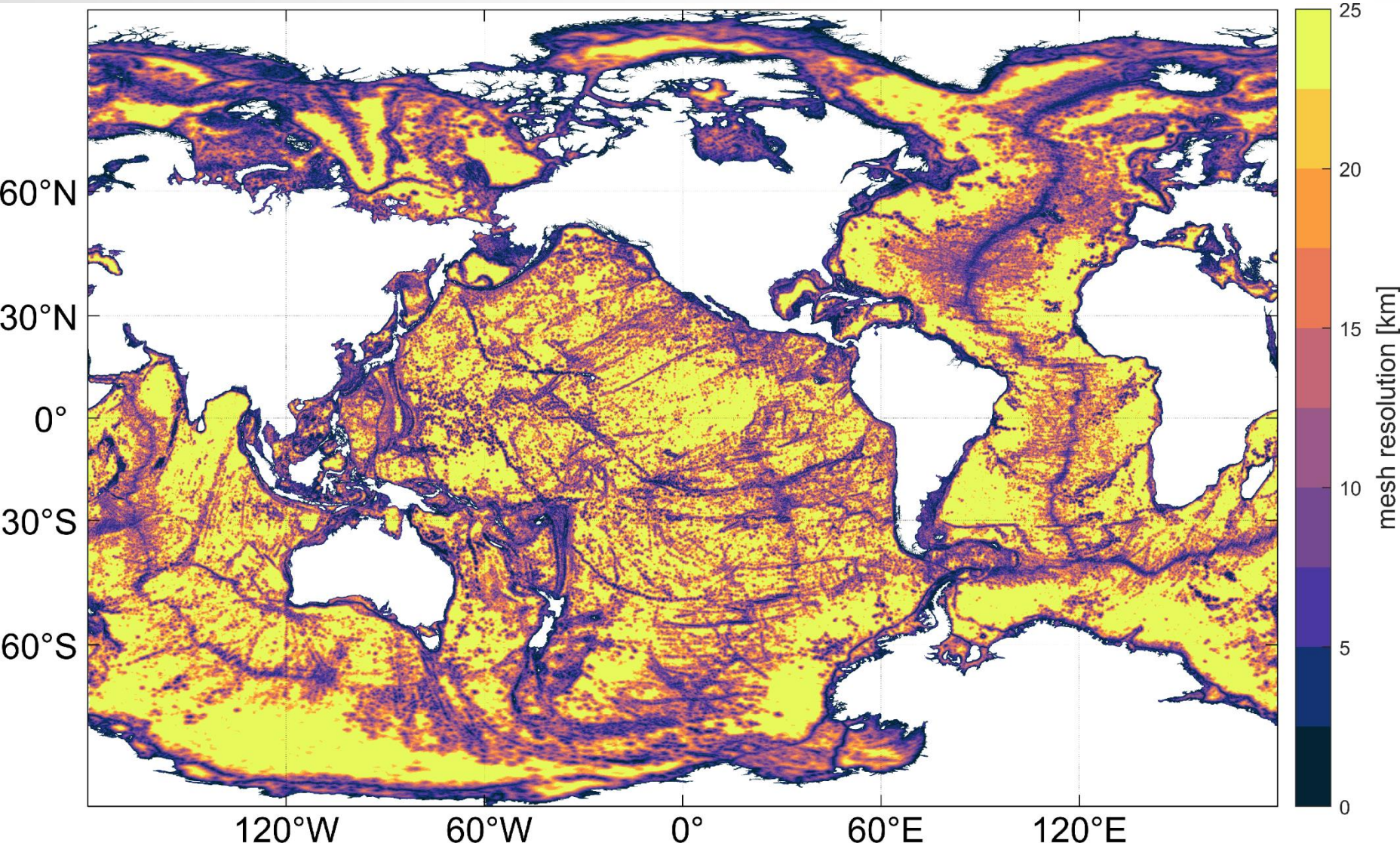


- Set **ICS** to a **negative** value, ex.:
ICS = -22
<https://wiki.adcirc.org/wiki/ICS>
- Supply a **fort.rotm** file, ex.:
znorth_in_spherical_coors
114.16991 0.77432
<https://wiki.adcirc.org/wiki/Fort.rotm>

- Allow the use of high mesh resolution at the physical pole.

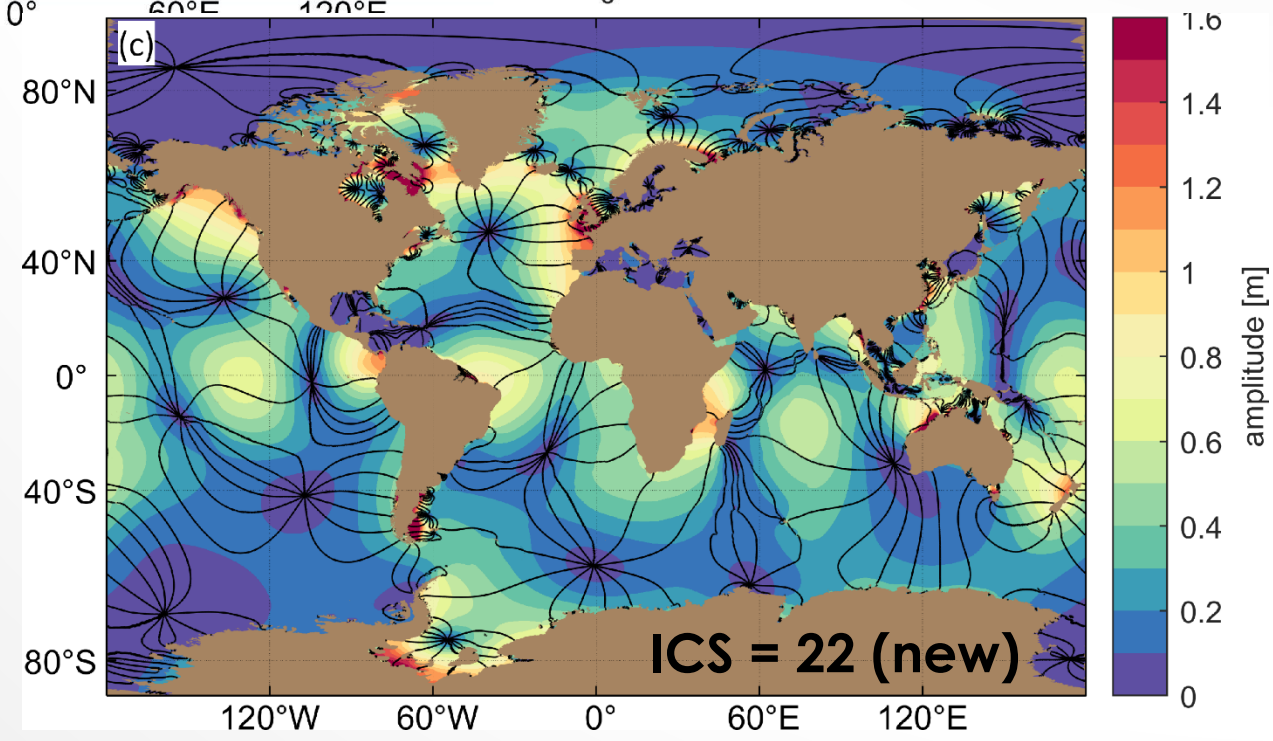
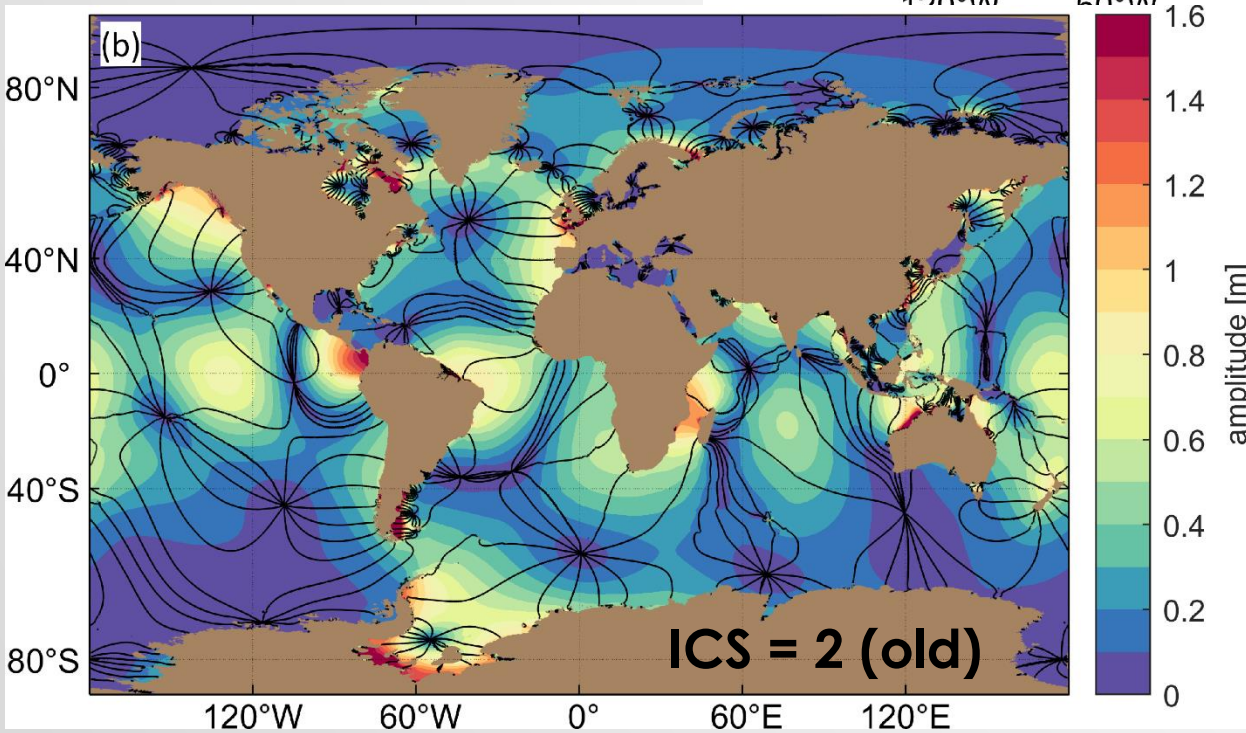
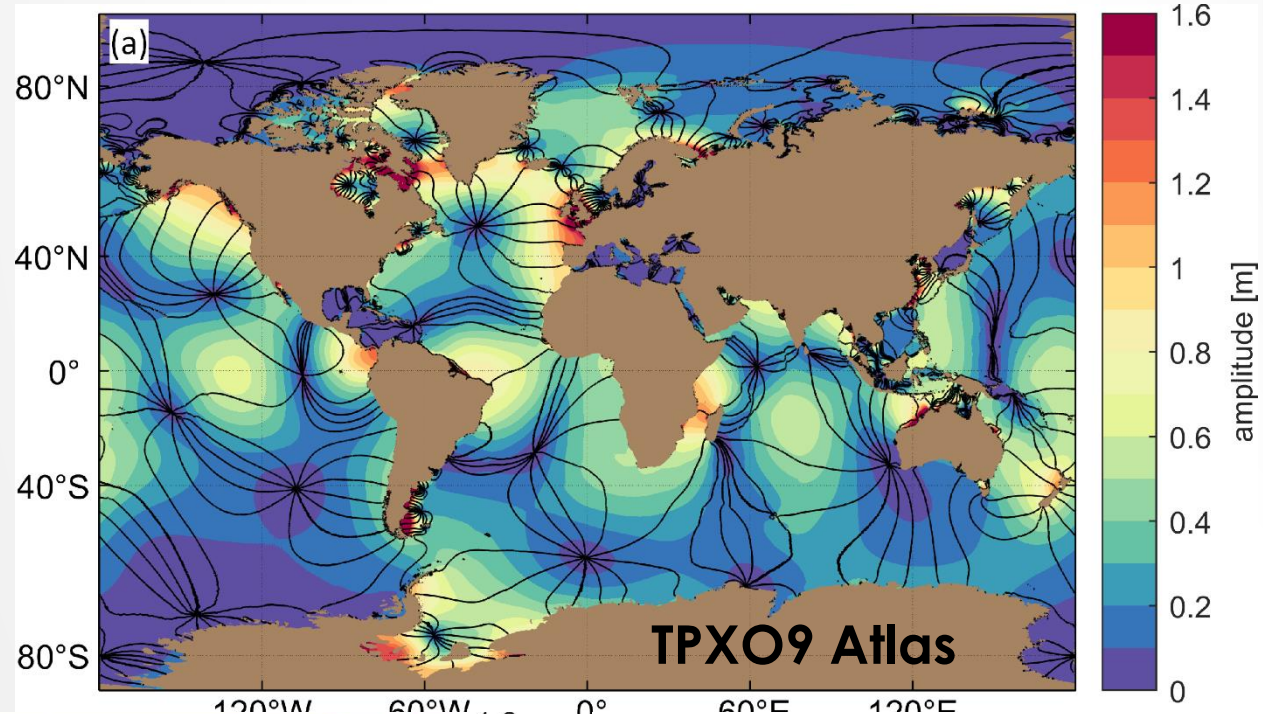
FIRST: TEST TIDES ON A HIGHLY RESOLVED MESH

Mesh rotated to have Greenland-Antarctica poles

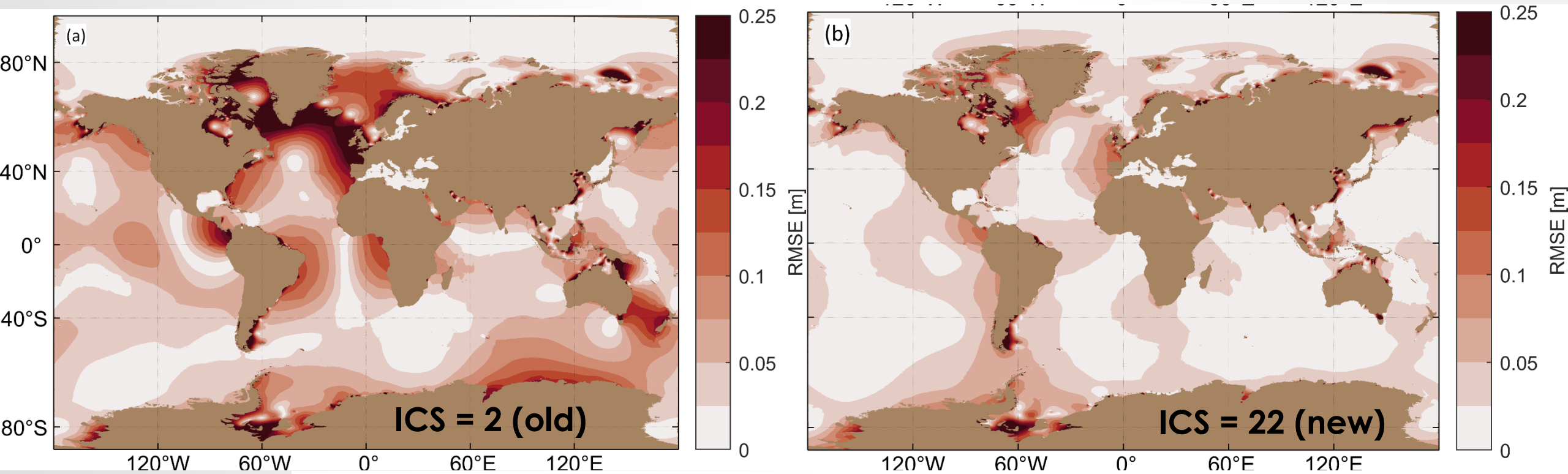


- Automatically generated by **[OceanMesh2D](#)**
[See: Example 7 Global.m](#)
- 6 million vertices
- 1.5 km to 25 km resolution
- Highly resolved along topographic gradients

M2 TIDES



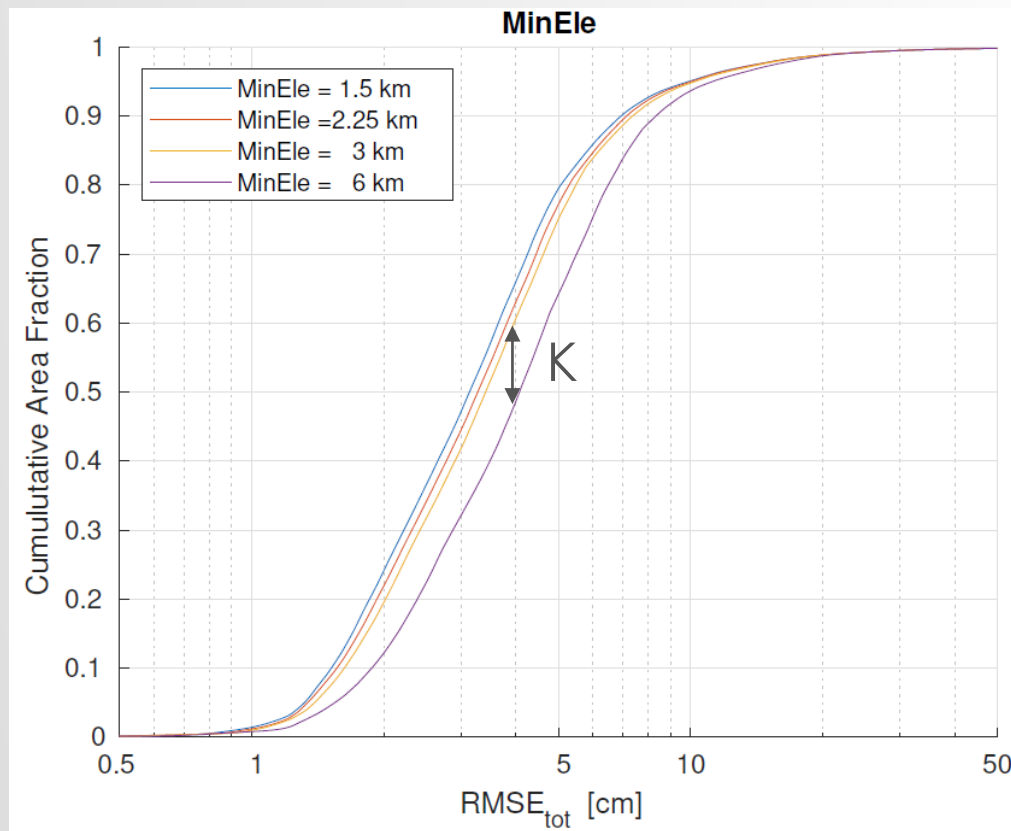
M2 RMSE



Model	$\overline{M_2 \text{ RMSE}_t}$ [cm]		$\overline{\text{RMSE}_{t tot}}$ [cm] [M2, N2, S2, K1, O1]	
	Deep water	Shallow water	Deep water	Shallow water
ADCIRC v54 ICS = 2 (old)	6.5	18.5	7.92	22.1
ADCIRC v55 ICS = 22 (new)	2.87	13.9	3.89	17.2 (Median = 6.67 cm)
Stammer et al. (2014)*	5.25-7.76	18.6-27.9	-	-
Ngodock et al. (2016)*	2.6-3.2	-	-	-

*: $\overline{\text{RMSE}_t}$ is computed against TPXO8-Atlas rather than TPXO9-Atlas.

SENSITIVITY TO MESH RESOLUTION

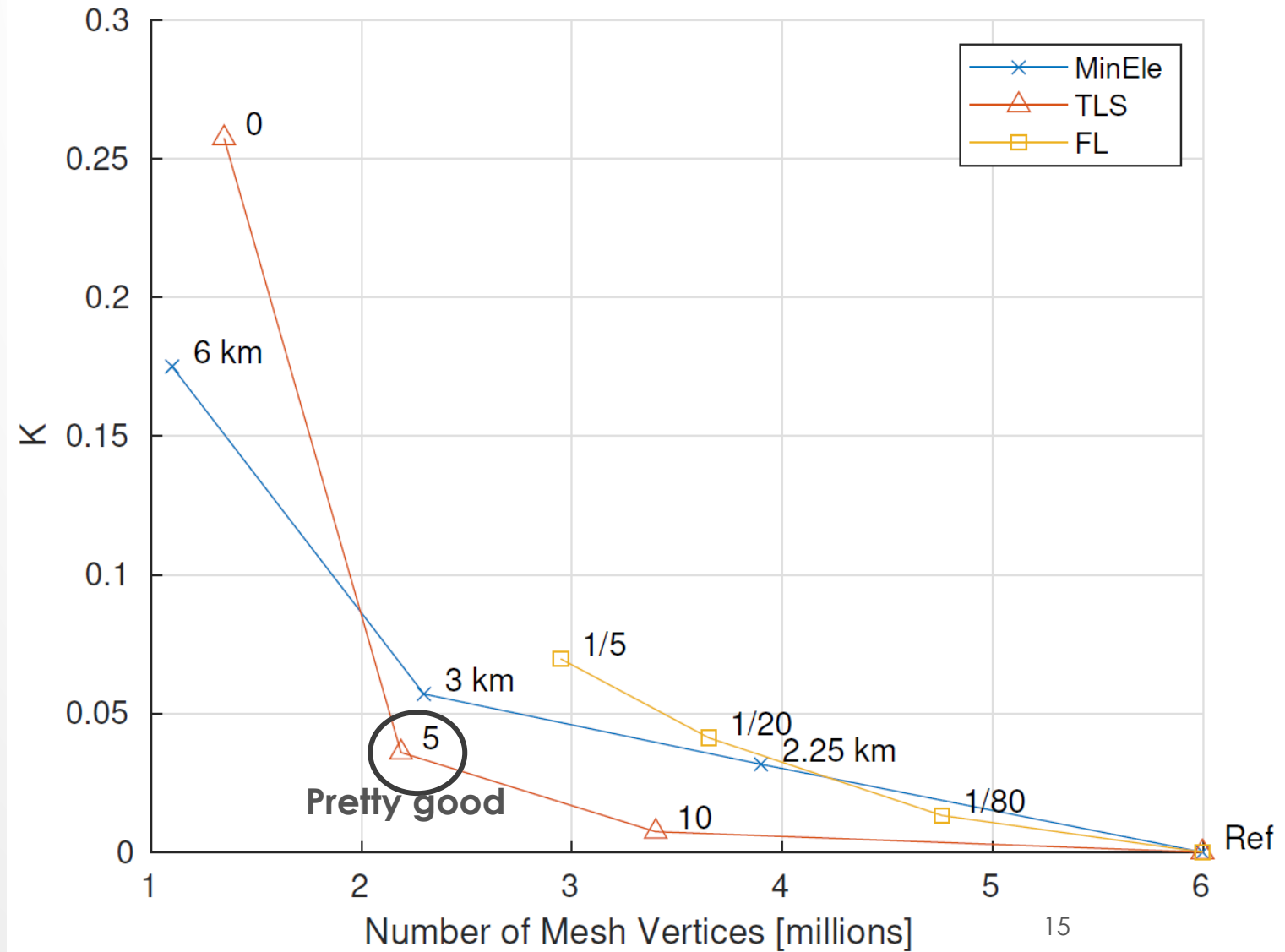


**edgex options in
OceanMesh2D**

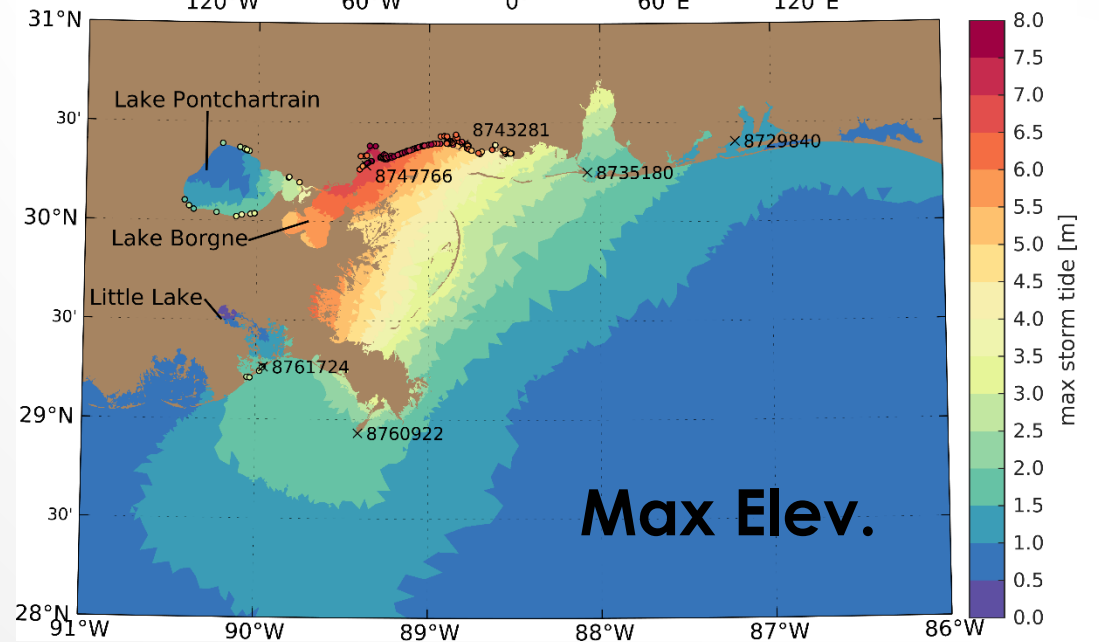
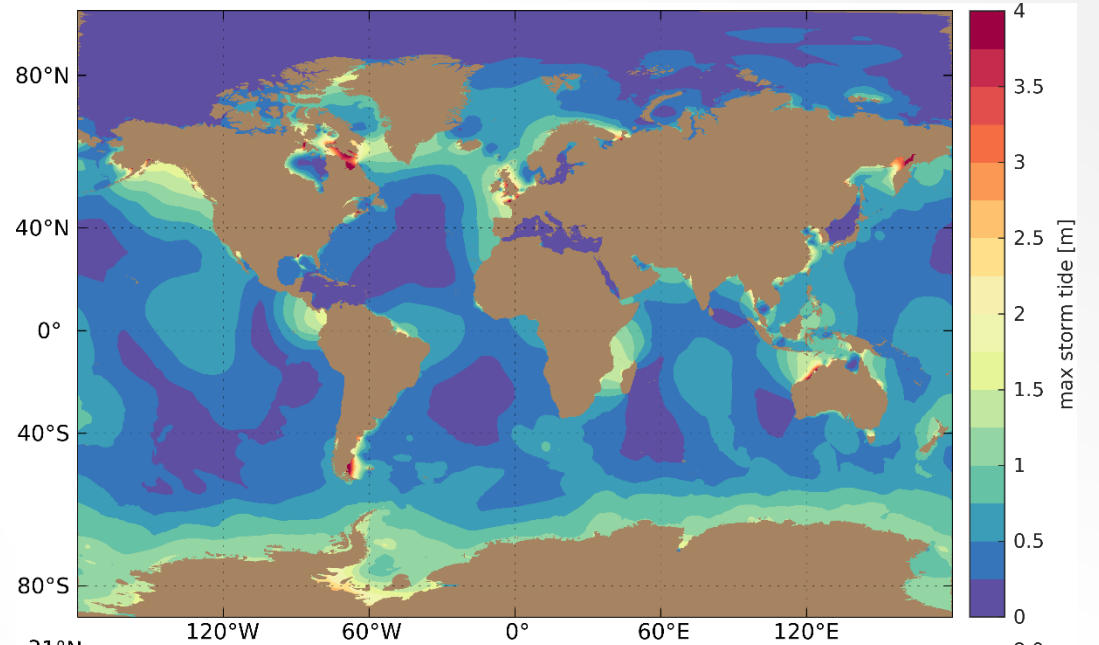
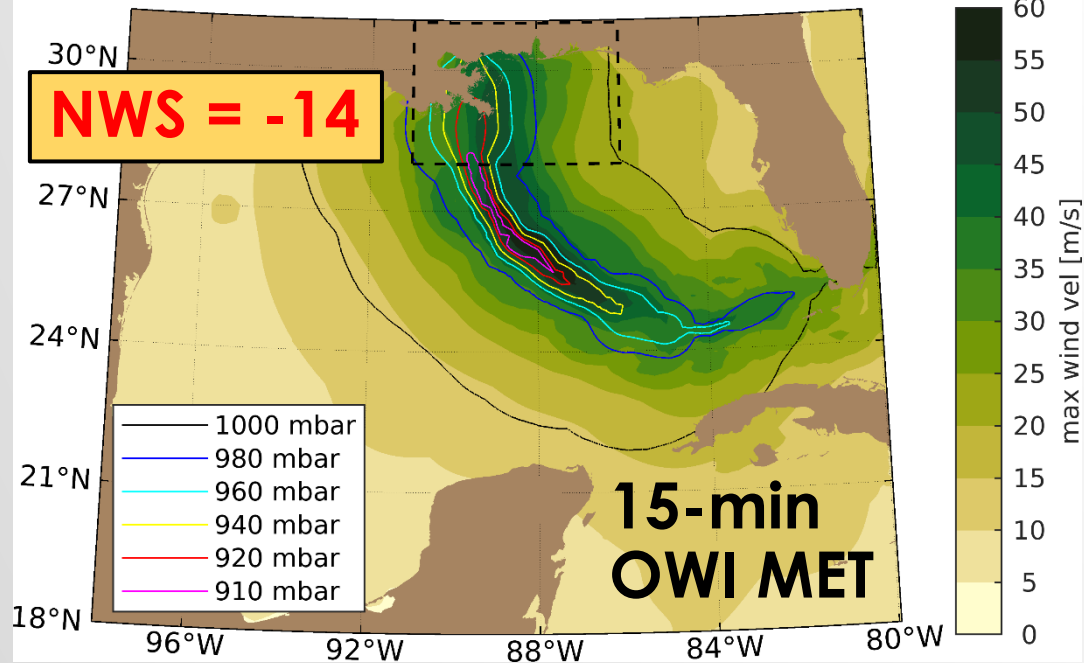
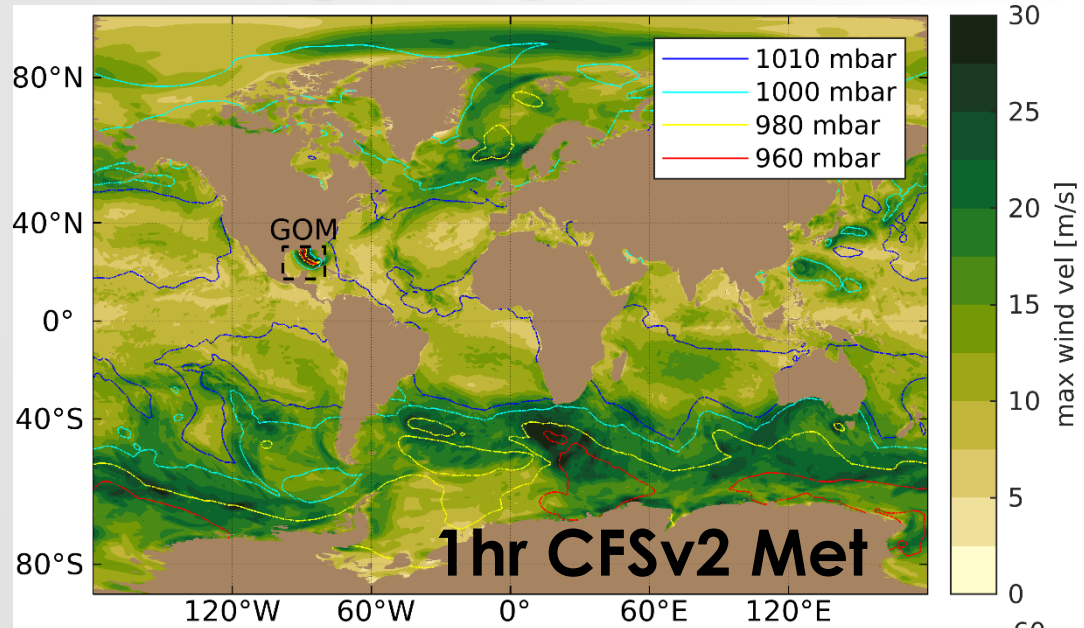
MinEle = 'h0'

TLS = 'slp'

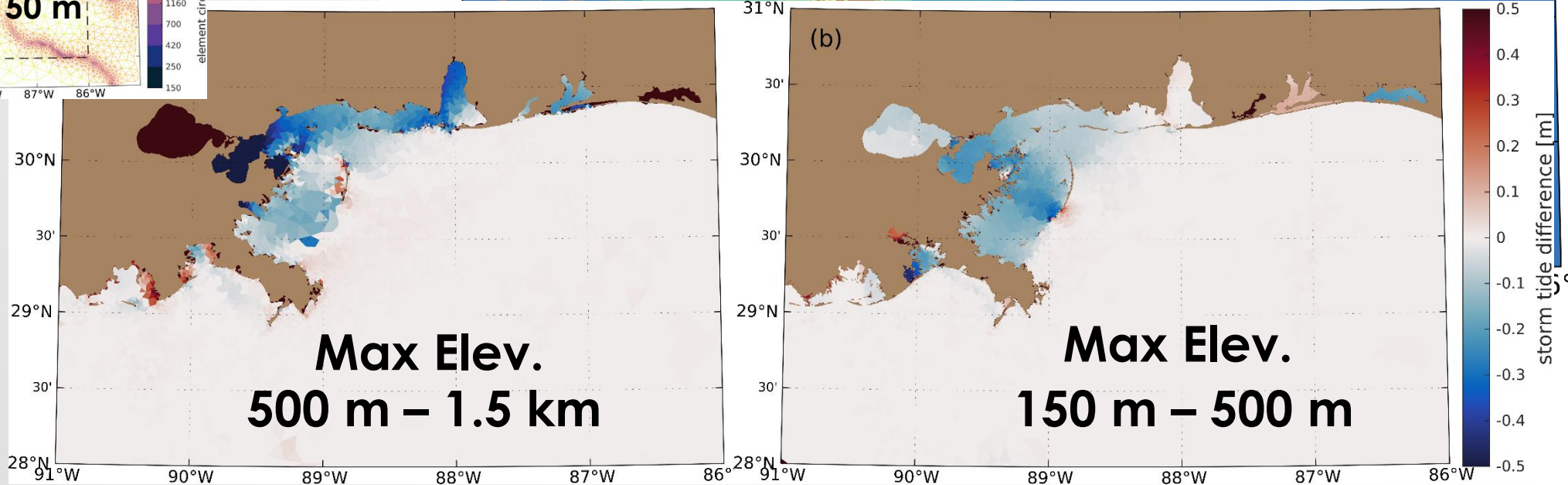
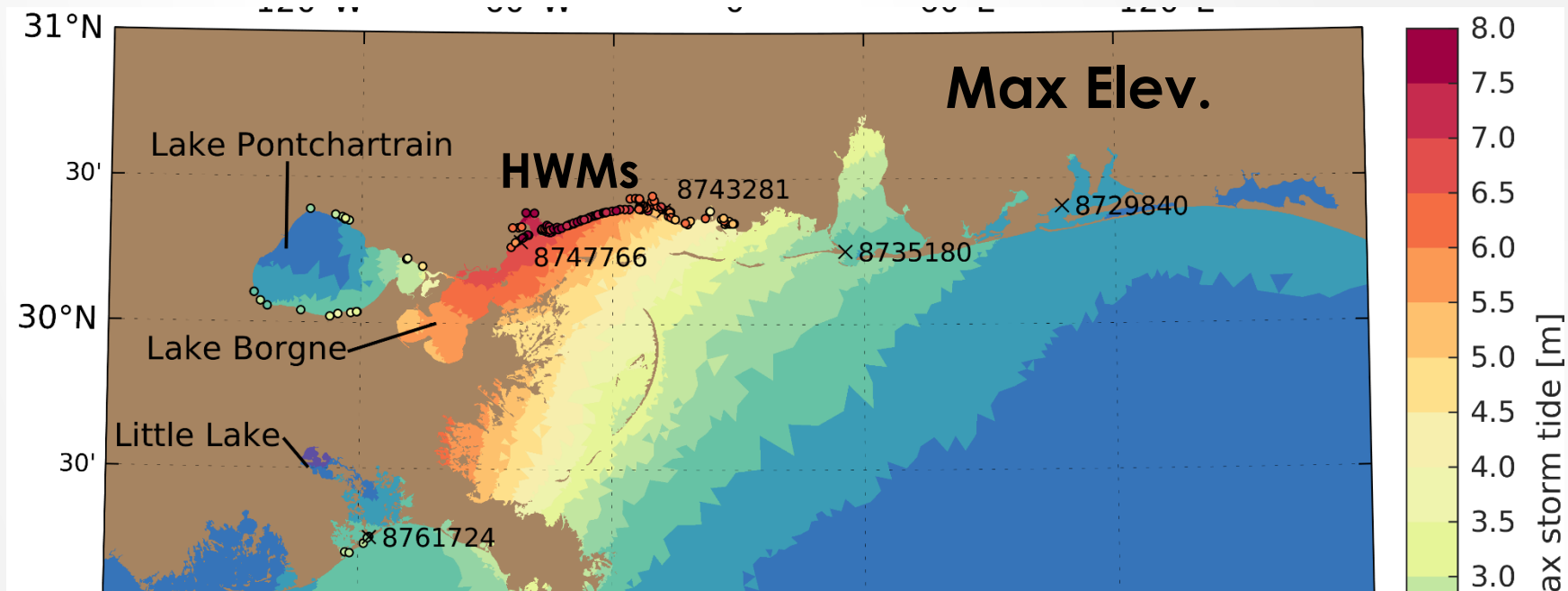
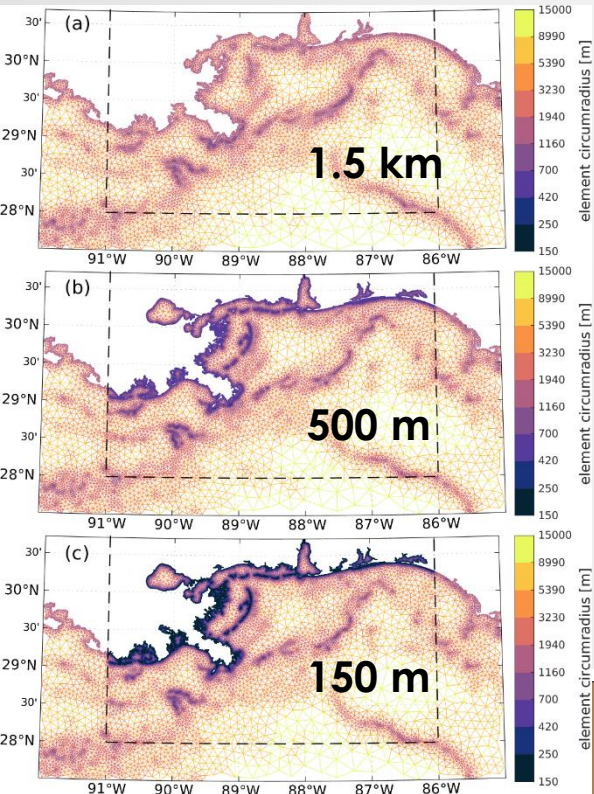
FL = 'fl'



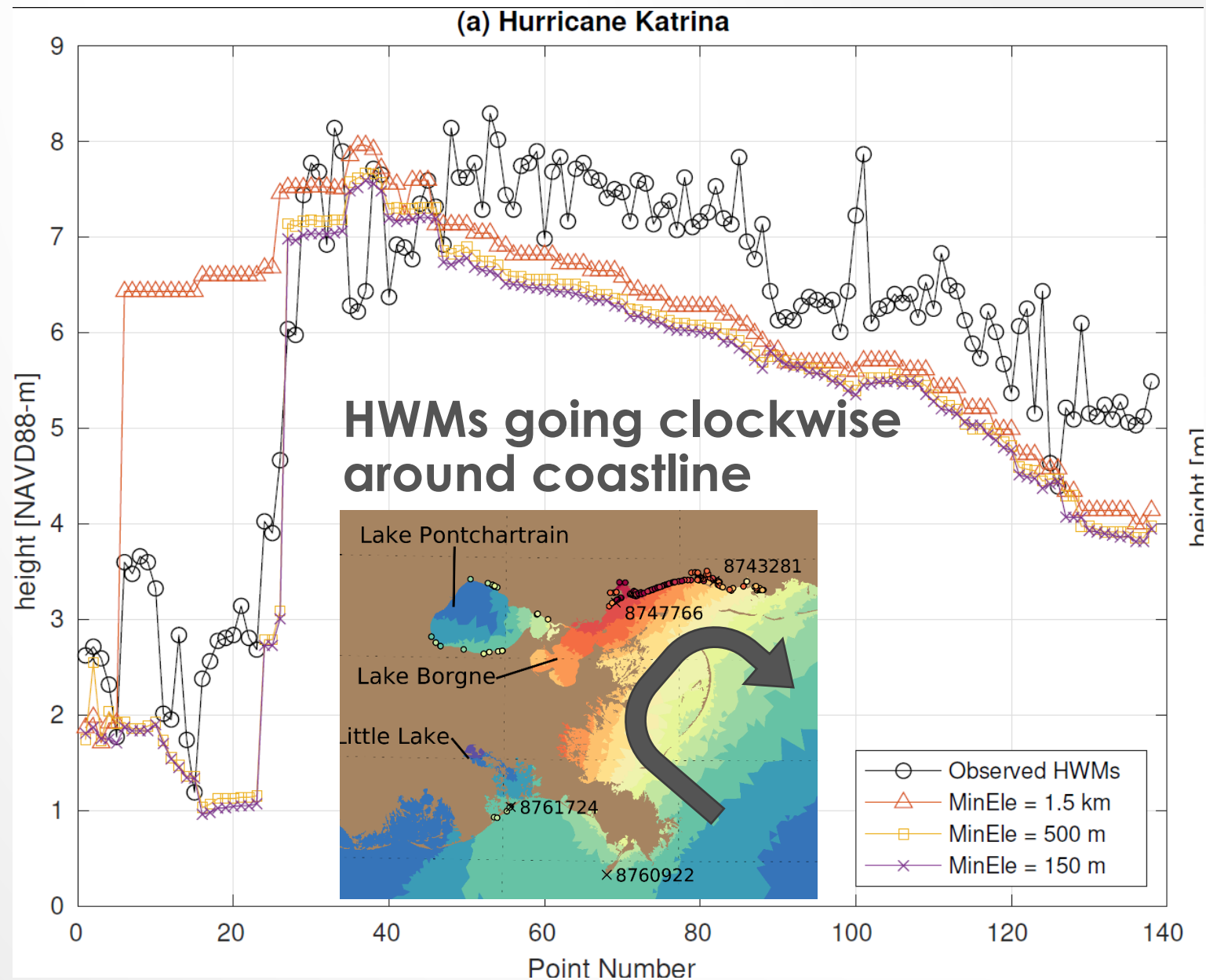
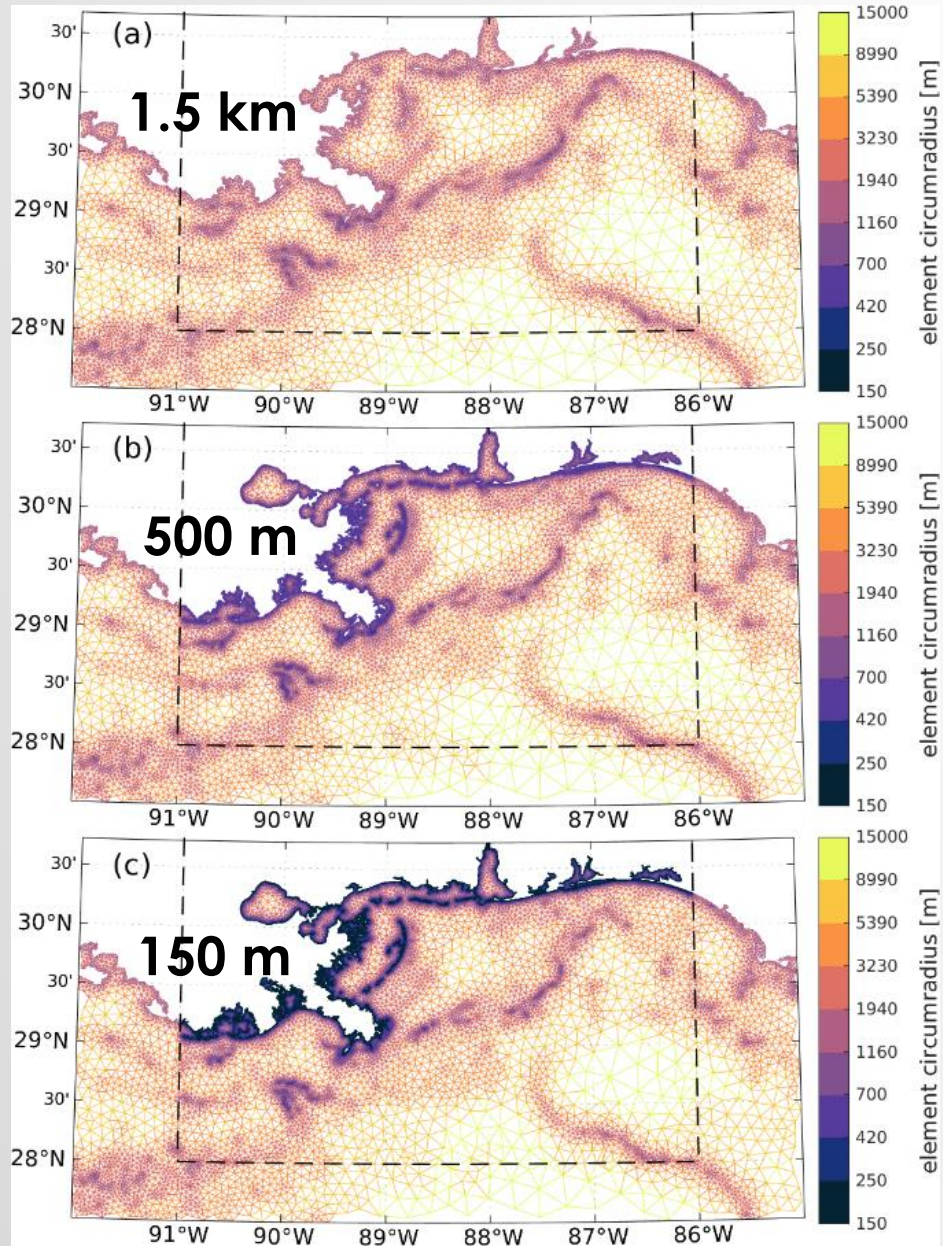
HURRICANE KATRINA ON GLOBAL MESH



SENSITIVITY TO LOCAL MESH REFINEMENT



SENSITIVITY TO LOCAL MESH REFINEMENT



2) IMPROVING STABILITY AND MASS CONSERVATION ISSUES

- What is the stability criteria for the GWCE? Let's check

1D linear case:

$$\begin{aligned} \frac{\zeta^{*s+1}}{\Delta t} \left(\frac{1}{\Delta t} + \frac{\tau_0}{2} \right) + \alpha_1 gh \frac{\partial^2 \zeta^{*s+1}}{\partial x^2} = \frac{\zeta^{*s}}{\Delta t} \left(\frac{1}{\Delta t} - \frac{\tau_0}{2} \right) \\ - (\alpha_1 + \alpha_2) gh \frac{\partial^2 \zeta^s}{\partial x^2} - \alpha_3 gh \frac{\partial^2 \zeta^{s-1}}{\partial x^2} + (\tau_0 - \tau_b^s) \frac{\partial U^s}{\partial x} \end{aligned}$$

$$U^{s+1} \left(\frac{1}{\Delta t} + \frac{\tau_b^s}{2} \right) = U^s \left(\frac{1}{\Delta t} - \frac{\tau_b^s}{2} \right) - \frac{g}{2} \left(\frac{\partial \zeta^s}{\partial x} + \frac{\partial \zeta^{s+1}}{\partial x} \right)$$

I follow method outlined in Kinnmark (1986) monograph based on Routh-Hurwitz criterion.

1) Make assumption on relationship between $\alpha_1, \alpha_2, \alpha_3$ (A00, B00, C00):

$\alpha_1 = \alpha_3 = \kappa, \text{ and } \alpha_2 = 1 - 2\kappa.$

e.g., **0.35, 0.30, 0.35**. Kinnmark (1986) analyzed this one

Equation	$Cr \rightarrow \infty$	$\tau_0 \rightarrow \infty$
p_0	$\kappa \geq 1/4$	$\kappa < 1/4 \text{ and } Cr^2 \leq \frac{1-m}{1-4\kappa}$
p_1	$\kappa \geq 1/4 \text{ and } \tau_0 \leq 4\kappa\tau_b$	$Cr^2 \leq 4(m/2 - 1)^2$
p_2	none	none
p_3	none	none
Δ_2	$\kappa \geq 1/4 \text{ and } \tau_0 \leq \tau_b$	$Cr^2 \leq 4(m/2 - 1)^2$

$$\tau_b = C_f \frac{U}{H}$$

$$\frac{\zeta^{*s+1}}{\Delta t} \left(\frac{1}{\Delta t} + \frac{\tau_0}{2} \right) + \alpha_1 gh \frac{\partial^2 \zeta^{*s+1}}{\partial x^2} = \frac{\zeta^{*s}}{\Delta t} \left(\frac{1}{\Delta t} - \frac{\tau_0}{2} \right) - (\alpha_1 + \alpha_2) gh \frac{\partial^2 \zeta^s}{\partial x^2} - \alpha_3 gh \frac{\partial^2 \zeta^{s-1}}{\partial x^2} + (\tau_0 - \tau_b^s) \frac{\partial U^s}{\partial x}$$

i.e., **TAU0** must be less than **linear bottom friction everywhere!** (which is impossible for quadratic friction as **U** \rightarrow **0**)

$$U^{s+1} \left(\frac{1}{\Delta t} + \frac{\tau_b^s}{2} \right) = U^s \left(\frac{1}{\Delta t} - \frac{\tau_b^s}{2} \right) - \frac{g}{2} \left(\frac{\partial \zeta^s}{\partial x} + \frac{\partial \zeta^{s+1}}{\partial x^{20}} \right)$$

Alternative relationship between $\alpha_1, \alpha_2, \alpha_3$ (A00, B00, C00):

$$\alpha_1 = \alpha_2 = \kappa, \text{ and } \alpha_3 = 1 - 2\kappa$$

Equation	$Cr \rightarrow \infty$	$\tau_0 \rightarrow \infty$
p_0	$\kappa \leq 1/2$	$\kappa > 1/2$ and $Cr^2 \leq \frac{1-m}{2\kappa-1}$
p_1	$1/3 \leq \kappa \leq 1/2$ and $\tau_0 \Delta t \leq 4(2-m)(3\kappa-1)$	$Cr^2 \leq 4(m/2-1)^2$
p_2	$\kappa \geq 1/3$	none
p_3	none	none
Δ_2	$\kappa \geq 1/3$	none

i.e., criteria between **TAU0** and **bottom friction vanishes!**

Example:

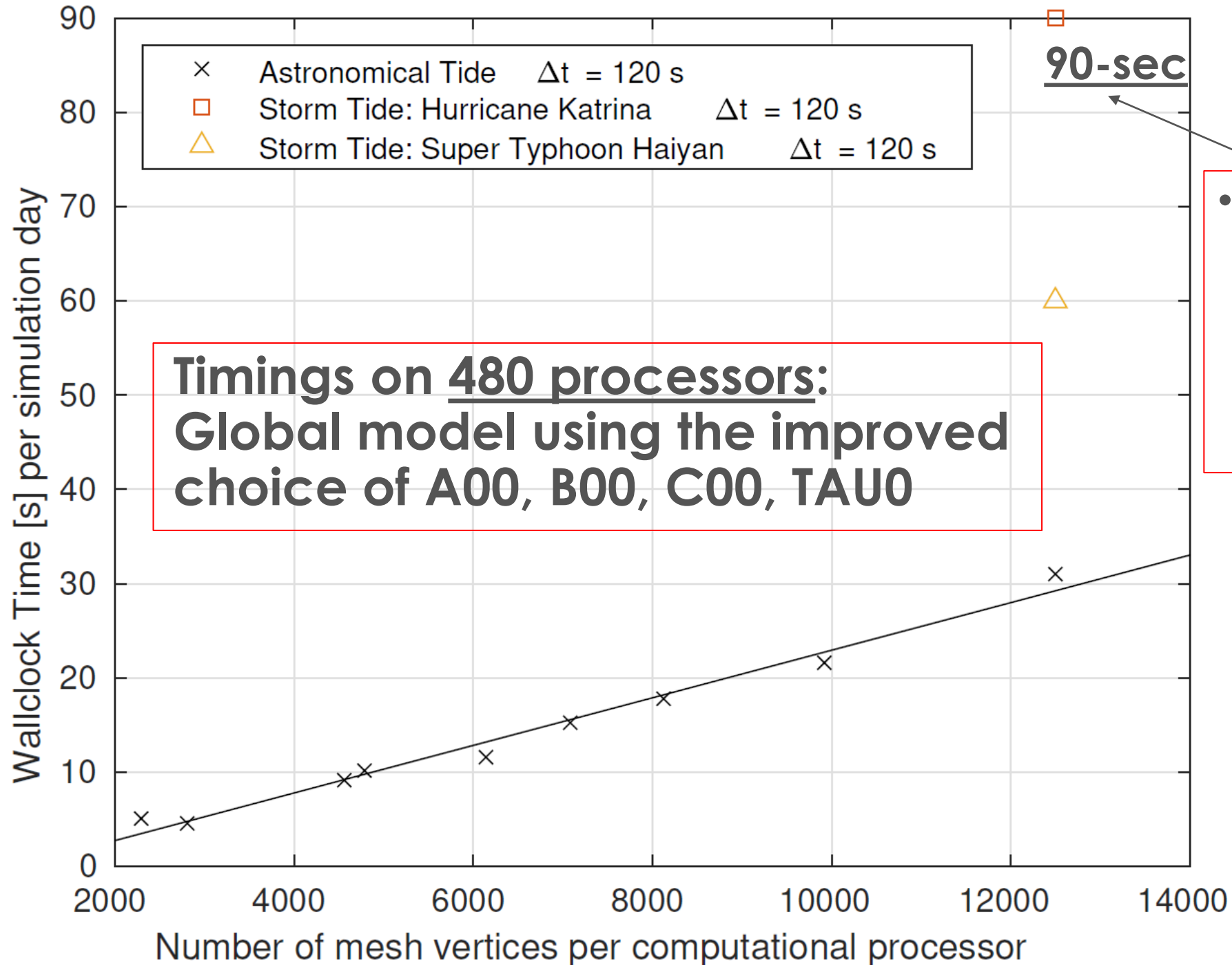
Select $\kappa = 0.5$, consistent mass-matrix: $m = 2/3$.

A00, B00, C00 = 0.5, 0.5, 0

TAU0 < 8/(3Δt)

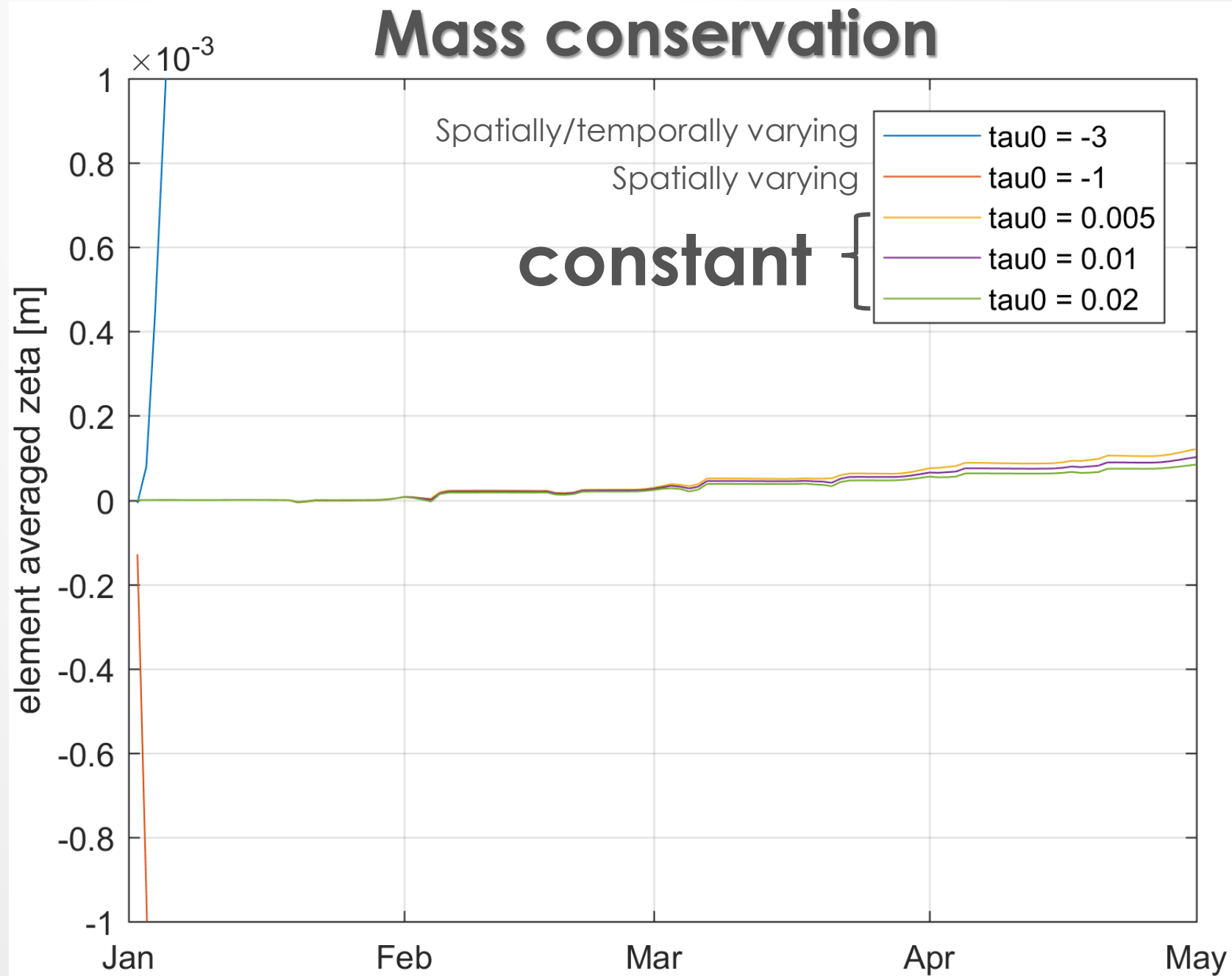
$$\begin{aligned} \frac{\zeta^{*s+1}}{\Delta t} \left(\frac{1}{\Delta t} + \frac{\tau_0}{2} \right) + \alpha_1 gh \frac{\partial^2 \zeta^{*s+1}}{\partial x^2} &= \frac{\zeta^{*s}}{\Delta t} \left(\frac{1}{\Delta t} - \frac{\tau_0}{2} \right) \\ &\quad - (\alpha_1 + \alpha_2) gh \frac{\partial^2 \zeta^s}{\partial x^2} - \alpha_3 gh \frac{\partial^2 \zeta^{s-1}}{\partial x^2} + (\tau_0 - \tau_b^s) \frac{\partial U^s}{\partial x} \end{aligned}$$

$$U^{s+1} \left(\frac{1}{\Delta t} + \frac{\tau_b^s}{2} \right) = U^s \left(\frac{1}{\Delta t} - \frac{\tau_b^s}{2} \right) - \frac{g}{2} \left(\frac{\partial \zeta^s}{\partial x} + \frac{\partial \zeta^{s+1}}{\partial x} \right)$$



- Old Hurricane Katrina simulations (Dietrich et al., 2011): **DT = 1 sec, 60-min** per simulation day on 480 processors

WHY WE SHOULD USE CONSTANT TAU0



SUMMARY

Pringle, W. J., et al. Global Storm Tide Modeling with ADCIRC v55 : Unstructured Mesh Design and Performance, will submit to Geosci. Model. Dev.

- 1) Extended **ADCIRC** to a **global** model
 - a. Tide solutions with **ICS = -22** clearly improved from old version (ICS = 2)
 - b. Resolution experiments show that global tide solutions decay beyond ~3 km min, and topographic slopes need to be well-resolved.
Local refinement generally decreases open ocean max. storm tide elevations
- 2) Improved **stability and mass-conservation** issues associated with **GWCE** formulation
 - a. Stability analysis shows that **A00, B00, C00 = 0.5, 0.5, 0** is non-Courant limited (in the linear sense) if **$\text{TAU0} < 8/(3\Delta t)$**

Hurricane Katrina Examples:

Global **1.5 km** mesh runs with $\Delta t = 120 \text{ s}$

Locally refined **500 m** runs with $\Delta t = 120 \text{ s}$

Locally refined **150 m** mesh runs with $\Delta t = 50 \text{ s}$

Also see:

<https://wiki.adcirc.org/wiki/IM>
for info on new option for fully
implicit gravity wave term
(**IM Digit 6 = 3**)

- b. From mass-conservation point of view we should use **constant TAU0**

GLobal Coastal Ocean Flood Forecasting System

An ADCIRC-based global storm tide modeling system providing real-time predictions for coastal flooding

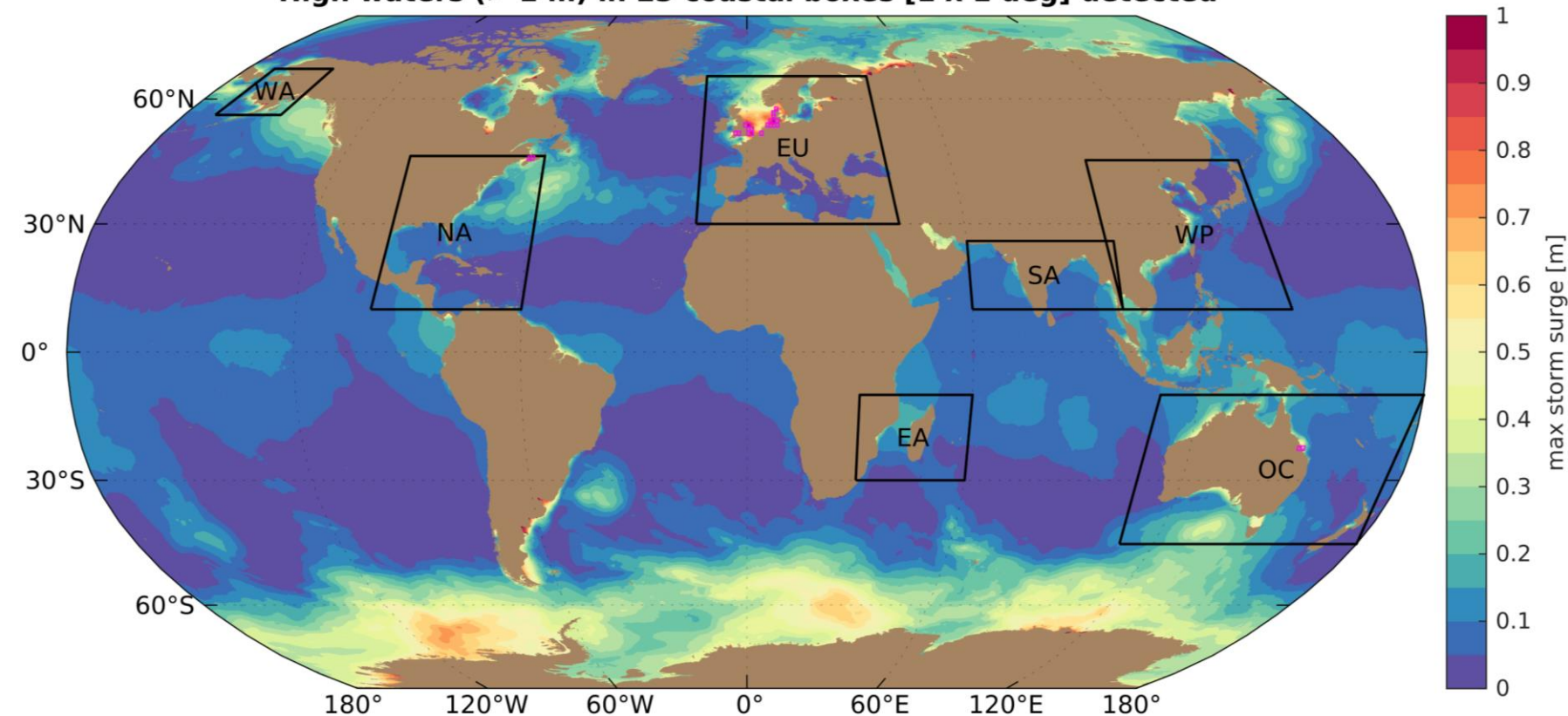
[View on GitHub](#)

<https://wpringle.github.io/GLOCOFFS/>

Max Storm Surge

2020-03-29 12:00 UTC to 2020-04-03 12:00 UTC

High waters (> 1 m) in 23 coastal boxes [1 x 1 deg] detected



**Continuous 5-day
forecasts every 6 hrs**

**Since April, 2019
1330+ commits**